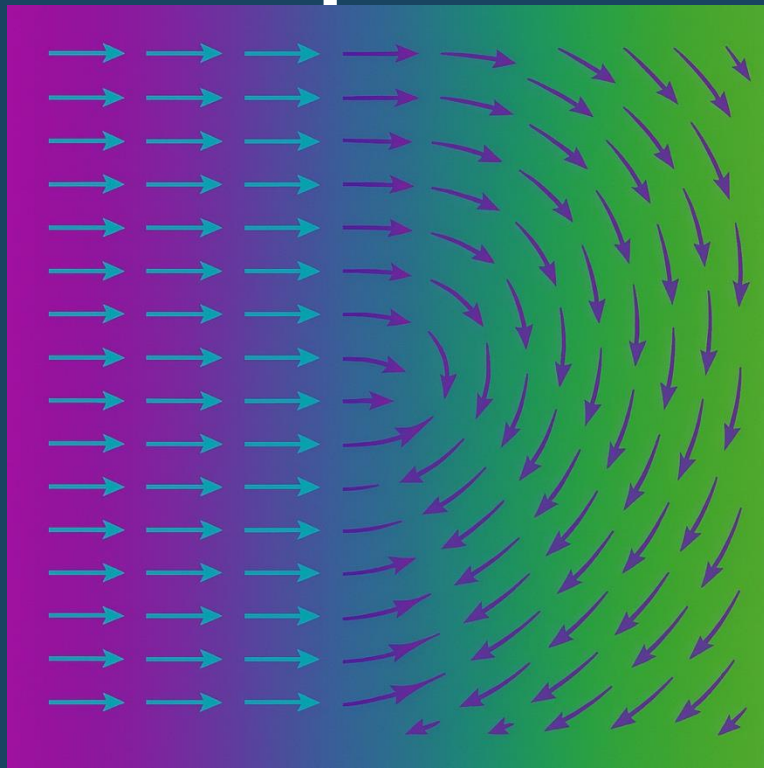


Ferrite now has true vector interpolations - what, why, and how?

Knut Andreas Meyer
Material and Computational Mechanics
Chalmers University of Technology



Outline

- **What** are true vector interpolations?
- **Why** are these useful?
- **How** are these implemented in Ferrite?
- **What** about boundary conditions?
- **Why** are these implemented in Ferrite?
- **How** should you use them?

What are vector interpolations

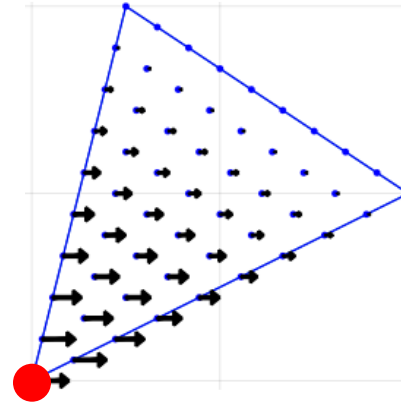
Current vectorized interpolations

- Based on scalar interpolations

$$N_1(\xi) = \begin{bmatrix} M_1(\xi) \\ 0 \end{bmatrix}, \quad N_3(\xi) = \begin{bmatrix} M_2(\xi) \\ 0 \end{bmatrix}, \quad N_5(\xi) = \begin{bmatrix} M_3(\xi) \\ 0 \end{bmatrix}$$

$$N_2(\xi) = \begin{bmatrix} 0 \\ M_1(\xi) \end{bmatrix}, \quad N_4(\xi) = \begin{bmatrix} 0 \\ M_2(\xi) \end{bmatrix}, \quad N_6(\xi) = \begin{bmatrix} 0 \\ M_3(\xi) \end{bmatrix}$$

- Vectors aligned with coordinate system
- Associated with algebraic nodes
- Inherits component-wise the continuity of the scalar interpolation



Lagrange

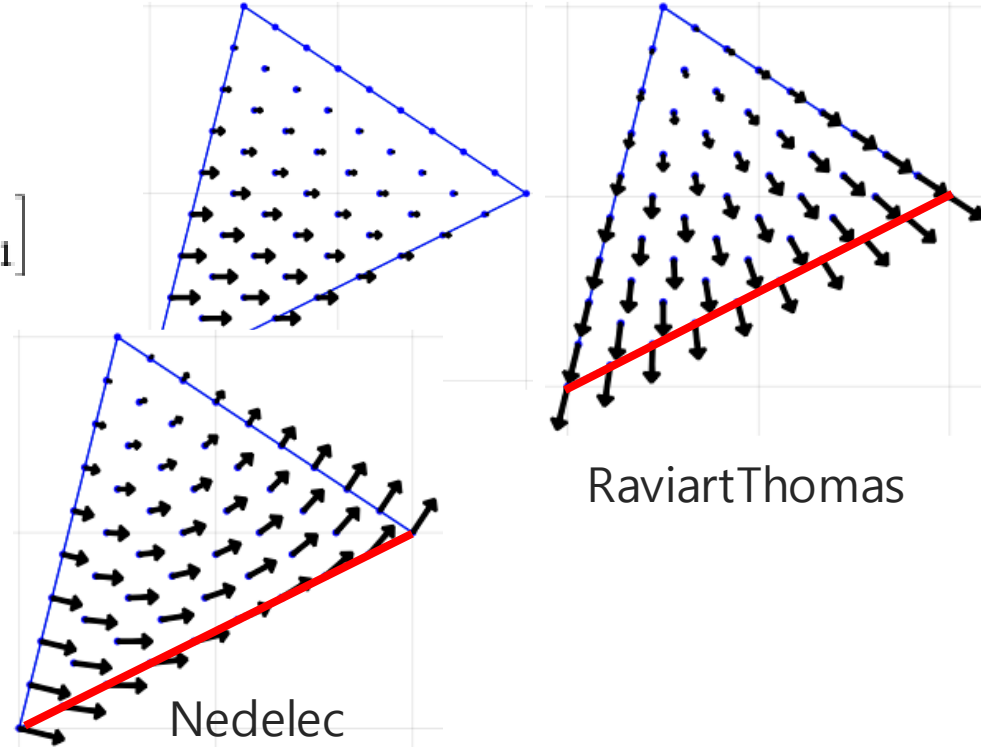
What are vector interpolations

True vector interpolations

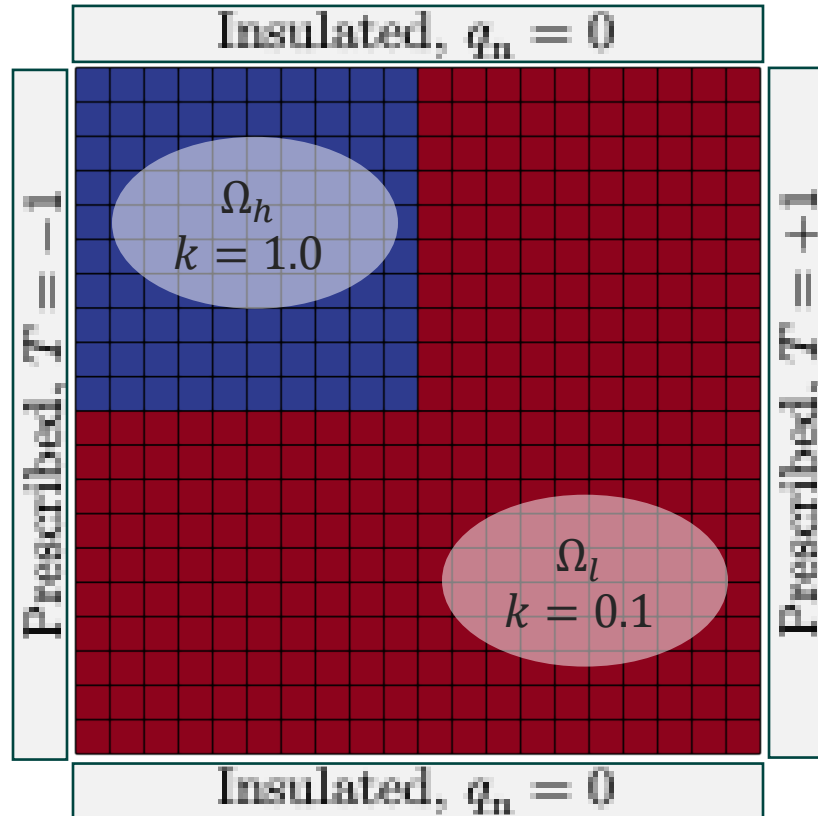
- Directly vector shape functions, e.g.
RaviartThomas{RefTriangle, 1}

$$N_1(\xi) = \begin{bmatrix} \xi_1 \\ \xi_2 \end{bmatrix}, \quad N_2(\xi) = \begin{bmatrix} \xi_1 - 1 \\ \xi_2 \end{bmatrix}, \quad N_3(\xi) = \begin{bmatrix} \xi_1 \\ \xi_2 - 1 \end{bmatrix}$$

- Vectors aligned with the cell
- Associated with different entities, e.g. edges and faces
- Allows normal/tangential continuity across element borders



Why? The heat equation

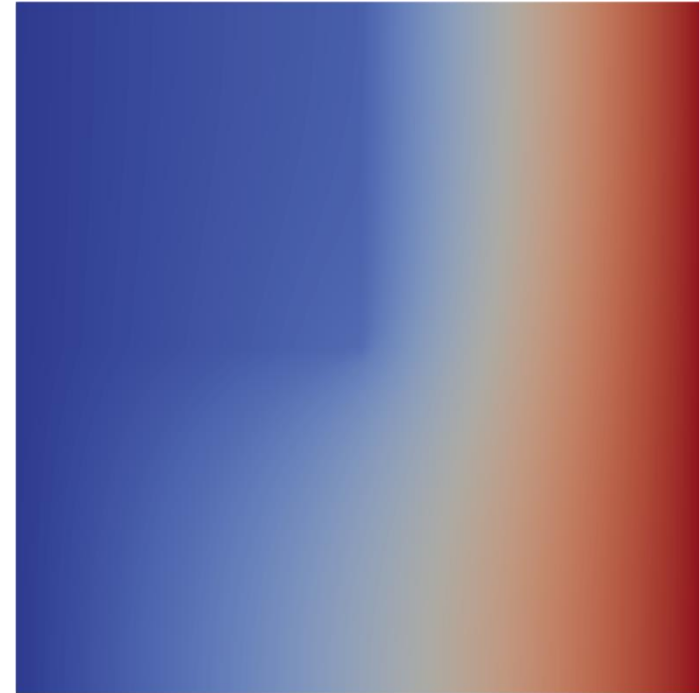
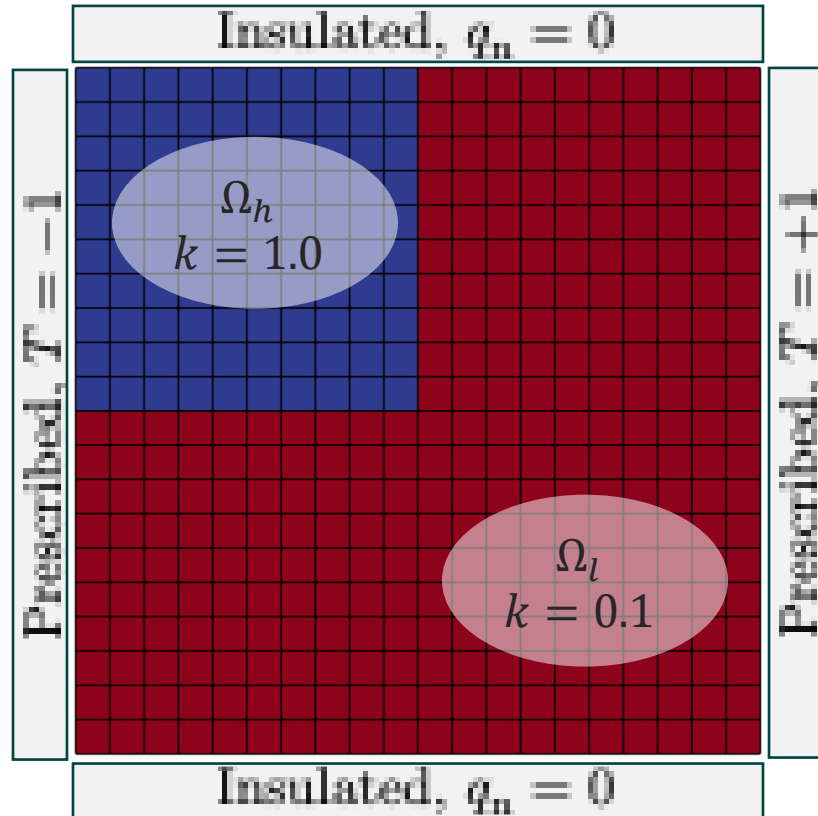


Energy conservation: $\mathbf{q}(\mathbf{x}) \cdot \nabla = 0$

Fourier's law: $\mathbf{q}(\mathbf{x}) = -k(\mathbf{x}) |\nabla T(\mathbf{x})|$

How will the heat flux field look like?

Why? The heat equation



Temperature



Why? The heat equation

How can we visualize the heat flux?

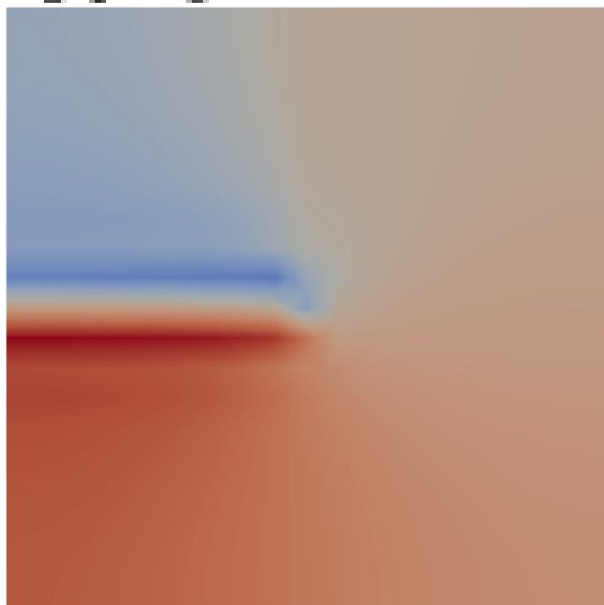
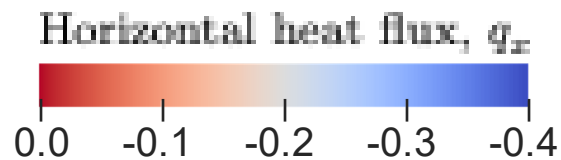
L_2 projection

- Data function, $f(\mathbf{x})$
- FE-approximation, $g(\mathbf{x}) = \sum_{i=1}^{N_{\text{dof}}} N_i(\mathbf{x}) a_i$
- minimize $e = \int_{\Omega} ||f(\mathbf{x}) - g(\mathbf{x})||^2 d\Omega$

$$\frac{1}{2} \frac{\partial e}{\partial a_i} = \underbrace{\int_{\Omega} N_i(\mathbf{x}) \cdot N_j(\mathbf{x}) d\Omega}_{K_{ij}} a_j - \underbrace{\int_{\Omega} N_i(\mathbf{x}) \cdot f(\mathbf{x}) d\Omega}_{f_i} = 0$$

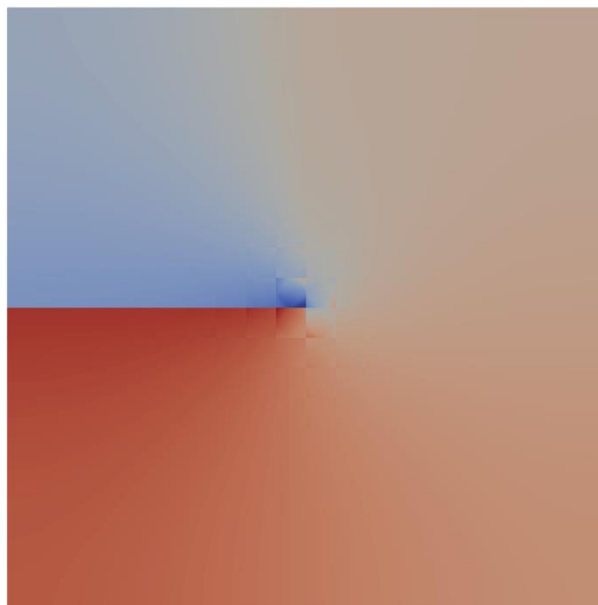
Why? The heat equation

L_2 projection results



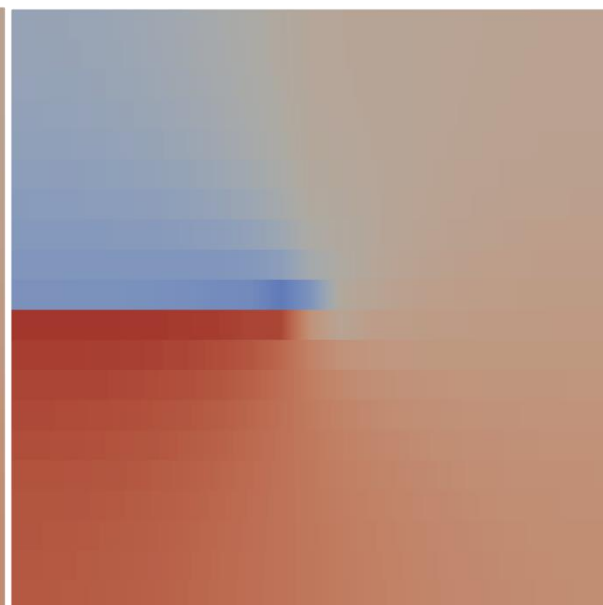
H^1 space

Lagrange{
RefQuadrilateral, 1}



L_2 space

DiscontinuousLagrange{
RefQuadrilateral, 1}



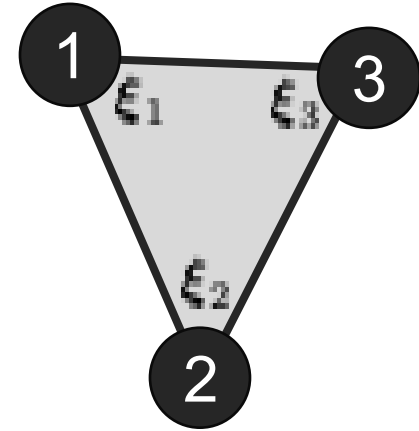
$H(\text{div})$ space

RaviartThomas{
RefQuadrilateral, 1}

How are these implemented?

Nodal interpolations

$$M_i(\xi_j) = \delta_{ij}$$



H(div) interpolations

$$\int_0^1 N_i(\xi = g_j(s)) f_j(s) \cdot \mathbf{n}_j \, ds = \delta_{ij} \quad \text{no sum on } j$$

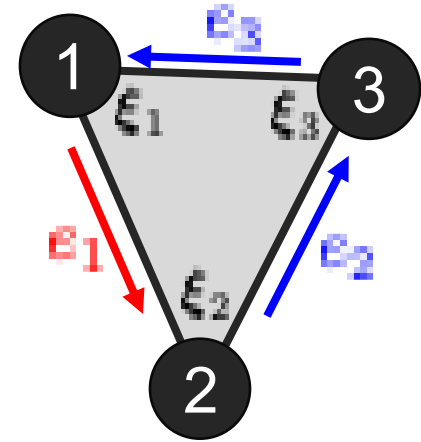
How are these implemented?

Vector interpolations, example: $H(\text{div})$

RaviartThomas{RefTriangle, 1}

$$\int_0^1 N_i(\xi_1 + vs) \cdot n_1 \|v\| \, ds = \begin{cases} 1 & i = 1 \\ 0 & i \neq 1 \end{cases}$$

$$v = \xi_2 - \xi_1$$



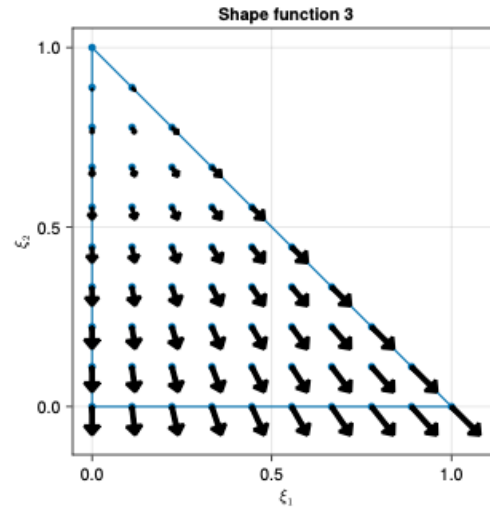
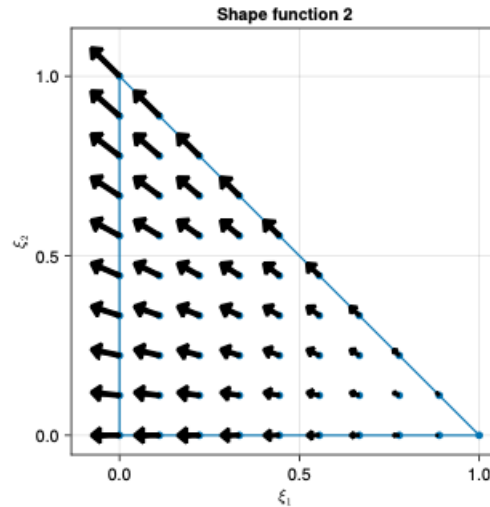
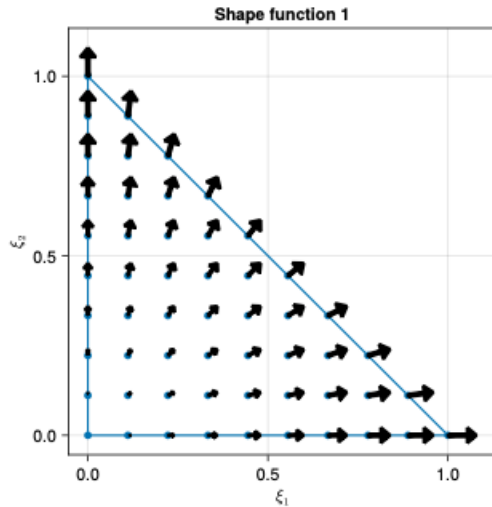
$H(\text{div})$ interpolations

$$\int_0^1 N_i(\xi = g_j(s)) f_j(s) \cdot n_j \, ds = \delta_{ij} \quad \text{no sum on } j$$

How are these implemented?

Vector interpolations, example: $H(\text{div})$

`RaviartThomas{RefTriangle, 1}`



How are these implemented?

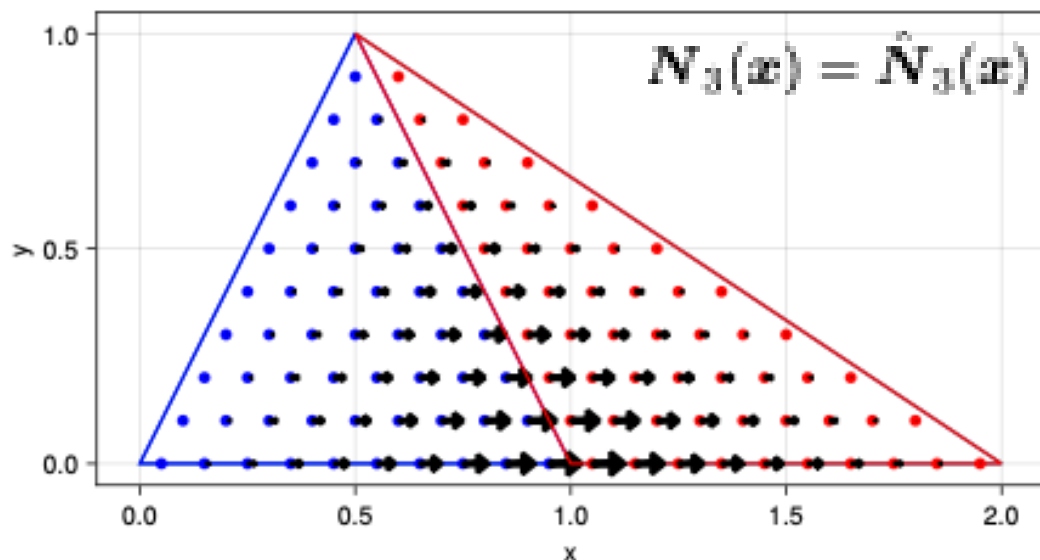
Vectorized interpolations, example:
 Lagrange{RefTriangle, 1}()^2

`Ferrite.IdentityMapping()`

$$N(\mathbf{x}) = \hat{N}(\xi(\mathbf{x}))$$

$$\frac{dN}{d\mathbf{x}} = \frac{d\hat{N}}{d\xi} \cdot \frac{d\xi}{d\mathbf{x}} = \frac{d\hat{N}}{d\xi} \cdot J^{-1}$$

$$J := \frac{d\mathbf{x}}{d\xi}$$



`N = shape_value(cellvalues, q_point, shapenr)`

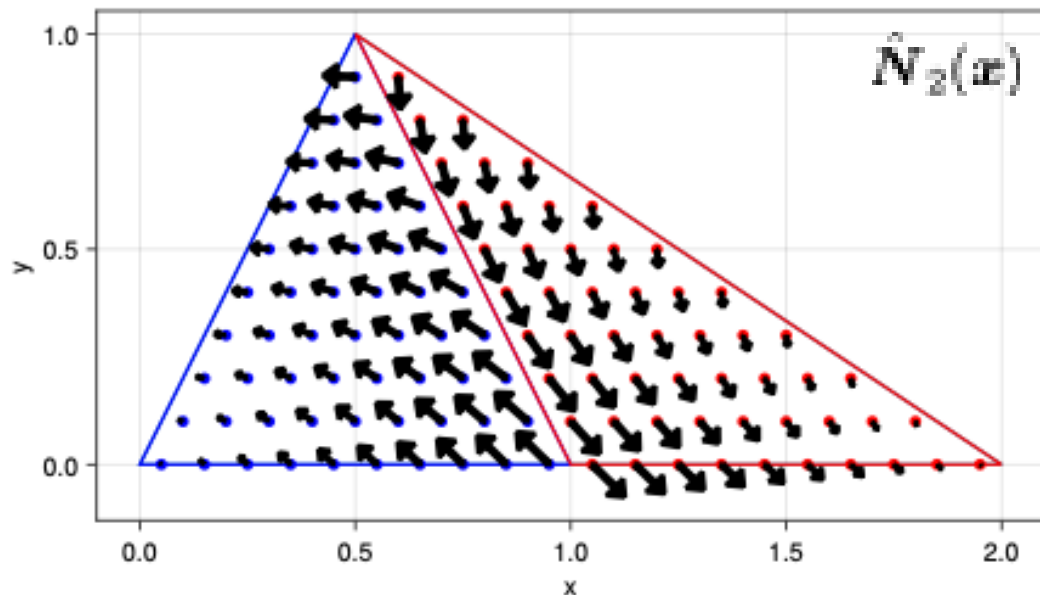
`N_hat = reference_shape_value(interpolation, local_coord, shapenr)`

How are these implemented?

Vector interpolations, example: H(div)
RaviartThomas{RefTriangle, 1}

`Ferrite.IdentityMapping()`

$$N(\mathbf{x}) = \hat{N}(\xi(\mathbf{x}))$$



`N = shape_value(cellvalues, q_point, shapenr)`

`N = reference_shape_value(interpolation, local_coord, shapenr)`

How are these implemented?

Vector interpolations, example: H(div)

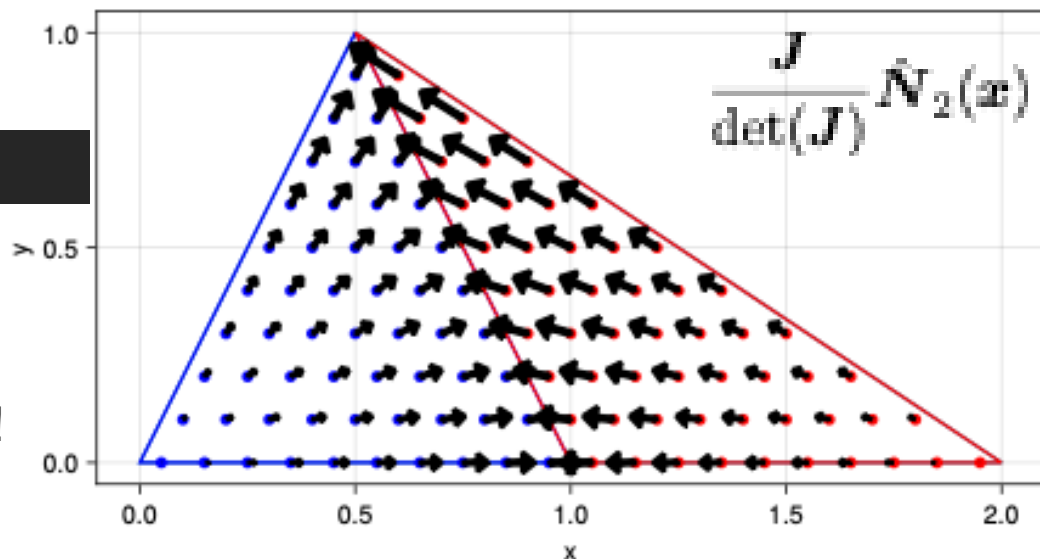
RaviartThomas{RefTriangle, 1}

`Ferrite.ContravariantPiolaMapping()`

$$\mathbf{N}(\mathbf{x}) = \frac{\mathbf{J}}{\det(\mathbf{J})} \hat{\mathbf{N}}(\boldsymbol{\xi}(\mathbf{x}))$$

Preserves the **normal** component!

But **sign** is wrong!



`N = shape_value(cellvalues, q_point, shapenr)`

`N̂ = reference_shape_value(interpolation, local_coord, shapenr)`

How are these implemented?

Vector interpolations, example: H(div)

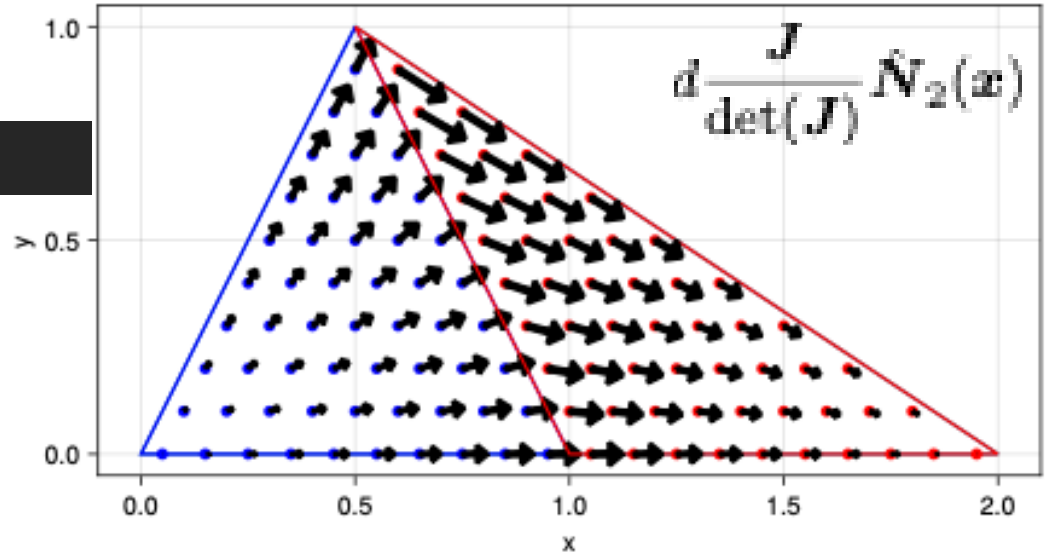
RaviartThomas{RefTriangle, 1}

`Ferrite.ContravariantPiolaMapping()`

$$\mathbf{N}(\mathbf{x}) = d \frac{\mathbf{J}}{\det(\mathbf{J})} \hat{\mathbf{N}}_2(\boldsymbol{\xi}(\mathbf{x}))$$

Introduce the facet direction:

$$d = \begin{cases} +1 & \text{positive facet} \\ -1 & \text{negative facet} \end{cases}$$



`N = shape_value(cellvalues, q_point, shapenr)`

`N = reference_shape_value(interpolation, local_coord, shapenr)`

How are these implemented?

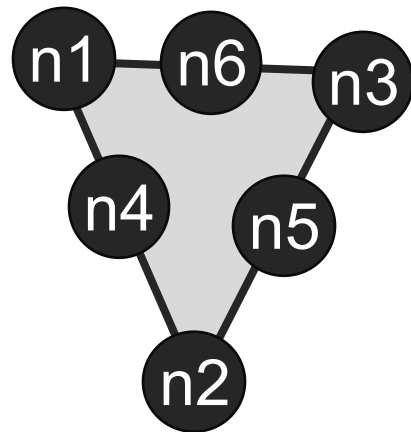
Facet direction

```
cell = QuadraticTriangle((n1, n2, n3, n4, n5, n6))
```

```
julia> facets(cell)  
((n1, n2), (n2, n3), (n3, n1))
```

```
# Consider facet 2: (n2, n3)
```

```
julia> d = n3 > n2 ? +1 : -1
```



This is why the **cell** is required to **reinit!** cellvalues with non-identity mapped interpolations

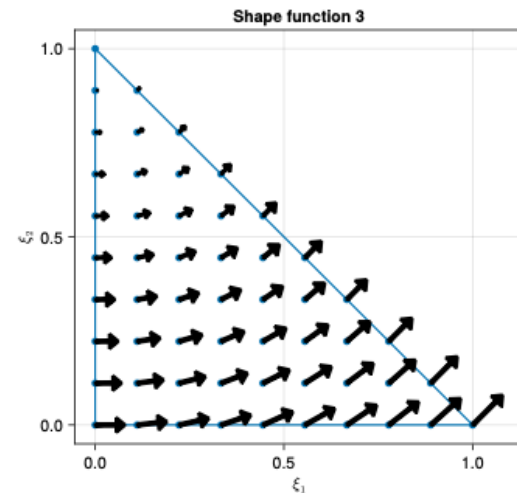
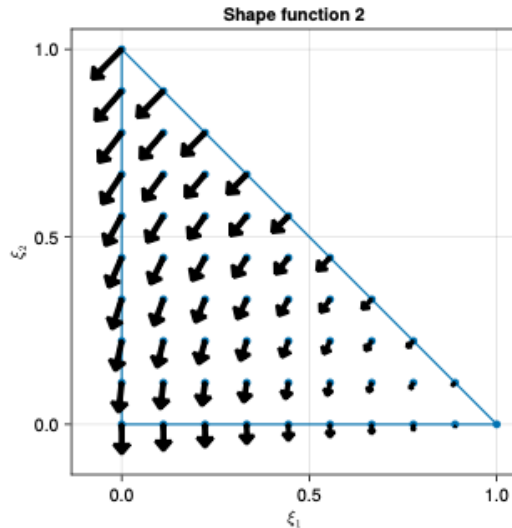
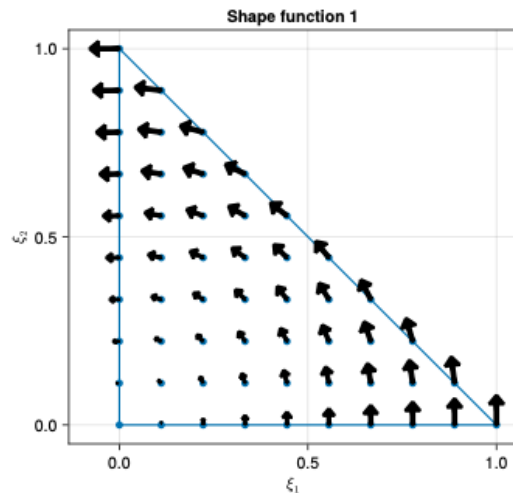
How are these implemented?

Vector interpolations, example: $H(\text{curl})$

Nedelec{RefTriangle, 1}

`Ferrite.CovariantPiolaMapping()`

Preserves the **tangential** component!



What about Boundary Conditions?

$$N_i(\mathbf{x})a_i = \mathbf{u}_{\text{prescribed}}(\mathbf{x}) \text{ on } \Gamma_D$$

$$N_i(\mathbf{x}_j) = \delta_{ij} \Rightarrow a_j = \mathbf{u}_{\text{prescribed}}(\mathbf{x}_j) \text{ for all } \mathbf{x}_j \text{ on } \Gamma_D$$

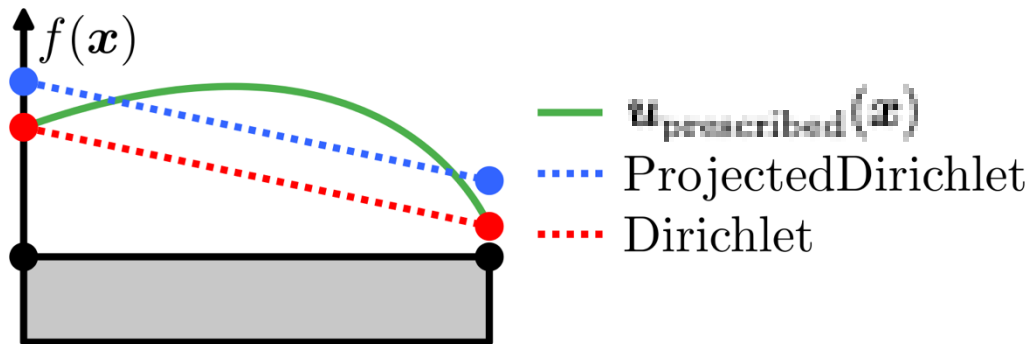
but $N_i(\mathbf{x}_j) \neq \delta_{ij}$ for true vector interpolations...

$$\text{minimize } \int_{\Gamma_D} [N_i(\mathbf{x})a_i - \mathbf{u}_{\text{prescribed}}(\mathbf{x})]^2 d\Gamma$$

```
pdbc = ProjectedDirichlet(  
  field_name::Symbol,  
  facets::Set{FacetIndex},  
  f::Function)
```

```
ch = ConstraintHandler(dh)  
add!(ch, pdbc)  
close!(ch)
```

What about Boundary Conditions?



$$a_j = u_{\text{prescribed}}(x_j)$$

$$\text{minimize } \int_{\Gamma_D} [N_i(x)a_i - u_{\text{prescribed}}(x)]^2 d\Gamma$$

```
pdbc = ProjectedDirichlet(
  field_name::Symbol,
  facets::Set{FacetIndex},
  u_prescribed::Function)
```

```
ch = ConstraintHandler(dh)
add!(ch, pdbc)
close!(ch)
```

Why are these implemented in Ferrite?

 Closed

Multiple cellvalues #674

 Closed

MultiCellValues #680

 Merged

Split values into geometric mapping and function values #764

 Open

CellMultiValues (new attempt) #872

 Merged

Implement $H(\text{div})$ and $H(\text{curl})$ interpolations #1145

```
ip = Lagrange{RefTriangle, 1}()  
cv = CellMultiValues(qr, ip, (u=ip^2, p=ip))
```

How should you use them?

- L2Projection (#1161)
<https://ferrite-fem.github.io/Ferrite.jl/previews/PR1161/howto/L2flux/>
- Heat equation: Mixed $H(\text{div})$ conforming formulation (#798)
https://ferrite-fem.github.io/Ferrite.jl/previews/PR798/tutorials/heat_equation_hdiv/
- Maxwell Discretizations: The Good, The Bad, and The Ugly (#798)
https://ferrite-fem.github.io/Ferrite.jl/previews/PR798/tutorials/maxwell_good_bad_ugly/

*Note that links are to previews
and will not work in the future*

Heat equation: Mixed $H(\text{div})$ conforming formulation

Theory

We start with the strong form of the heat equation: Find the temperature, $u(\mathbf{x})$, and heat flux, $\mathbf{q}(\mathbf{x})$, such that

$$\begin{aligned}\nabla \cdot \mathbf{q} &= h(\mathbf{x}), & \text{in } \Omega \\ \mathbf{q}(\mathbf{x}) &= -k \nabla u(\mathbf{x}), & \text{in } \Omega \\ \mathbf{q}(\mathbf{x}) \cdot \mathbf{n}(\mathbf{x}) &= q_n, & \text{on } \Gamma_N \\ u(\mathbf{x}) &= u_D, & \text{on } \Gamma_D\end{aligned}$$

From this strong form, we can formulate the weak form as a mixed formulation. Find $u \in \mathbb{U}$ and $\mathbf{q} \in \mathbb{Q}$ such that

$$\begin{aligned}\int_{\Omega} \delta u [\nabla \cdot \mathbf{q}] \, d\Omega &= \int_{\Omega} \delta u h \, d\Omega, \quad \forall \delta u \in \delta \mathbb{U} \\ \int_{\Omega} \delta \mathbf{q} \cdot \mathbf{q} \, d\Omega - \int_{\Omega} [\nabla \cdot \delta \mathbf{q}] k u \, d\Omega &= - \int_{\Gamma} \delta \mathbf{q} \cdot \mathbf{n} k u \, d\Gamma, \quad \forall \delta \mathbf{q} \in \delta \mathbb{Q}\end{aligned}$$

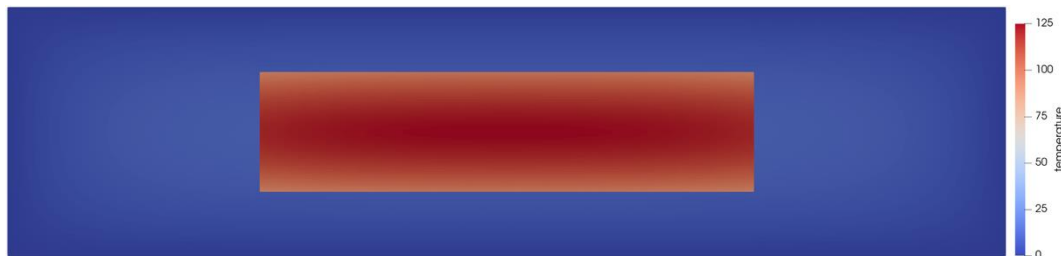
where we have the function spaces,

$$\begin{aligned}\mathbb{U} &= \delta \mathbb{U} = L^2 \\ \mathbb{Q} &= \{\mathbf{q} \in H(\text{div}) \text{ such that } \mathbf{q} \cdot \mathbf{n} = q_n \text{ on } \Gamma_D\} \\ \delta \mathbb{Q} &= \{\mathbf{q} \in H(\text{div}) \text{ such that } \mathbf{q} \cdot \mathbf{n} = 0 \text{ on } \Gamma_D\}\end{aligned}$$

A stable choice of finite element spaces for this problem on grid with triangles is using

- `DiscontinuousLagrange{RefTriangle, k-1}` for approximating L^2
- `BrezziDouglasMarini{RefTriangle, k}` for approximating $H(\text{div})$

Heat equation: Mixed $H(\text{div})$ conforming formulation



- Expected since temperature is discontinuous (L^2)
- Boundary flux “exact”

Theory

We start with the strong form of the heat equation: Find the temperature, $u(\mathbf{x})$, and heat flux, $\mathbf{q}(\mathbf{x})$, such that

$$\begin{aligned}\nabla \cdot \mathbf{q} &= h(\mathbf{x}), & \text{in } \Omega \\ \mathbf{q}(\mathbf{x}) &= -k \nabla u(\mathbf{x}), & \text{in } \Omega \\ \mathbf{q}(\mathbf{x}) \cdot \mathbf{n}(\mathbf{x}) &= q_n, & \text{on } \Gamma_N \\ u(\mathbf{x}) &= u_D, & \text{on } \Gamma_D\end{aligned}$$

From this strong form, we can formulate the weak form as a mixed formulation. Find $u \in \mathbb{U}$ and $\mathbf{q} \in \mathbb{Q}$ such that

$$\begin{aligned}\int_{\Omega} \delta u [\nabla \cdot \mathbf{q}] \, d\Omega &= \int_{\Omega} \delta u h \, d\Omega, \quad \forall \delta u \in \delta \mathbb{U} \\ \int_{\Omega} \delta \mathbf{q} \cdot \mathbf{q} \, d\Omega - \int_{\Omega} [\nabla \cdot \delta \mathbf{q}] k u \, d\Omega &= - \int_{\Gamma} \delta \mathbf{q} \cdot \mathbf{n} k u \, d\Gamma, \quad \forall \delta \mathbf{q} \in \delta \mathbb{Q}\end{aligned}$$

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A stable choice of finite element spaces for this problem on grid with triangles is using

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Maxwell Discretizations: The Good, The Bad, and The Ugly

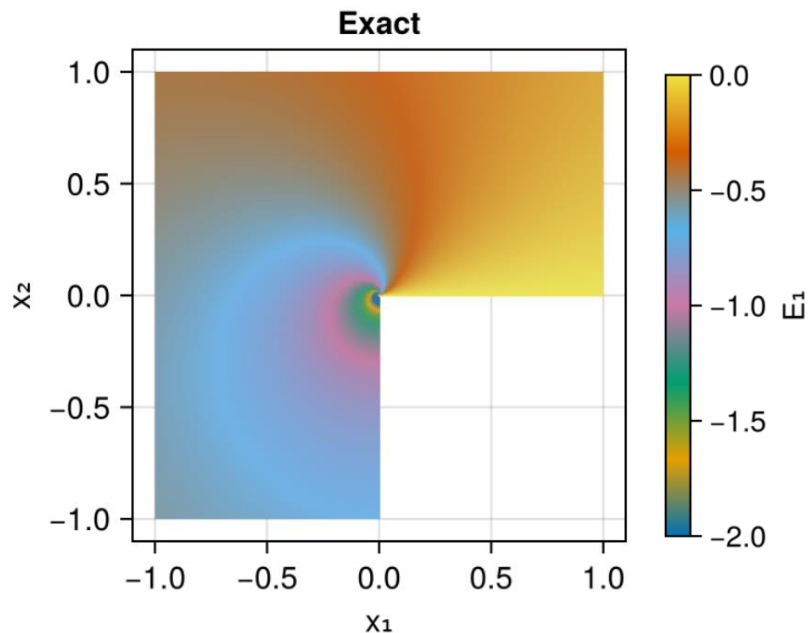
$$\text{curl}(\text{curl}(\mathbf{E})) = 0 \quad \text{in } \Omega$$

$$\text{div}(\mathbf{E}) = 0 \quad \text{in } \Omega$$

$$\mathbf{E} \cdot \mathbf{t} = g \quad \text{on } \Gamma$$

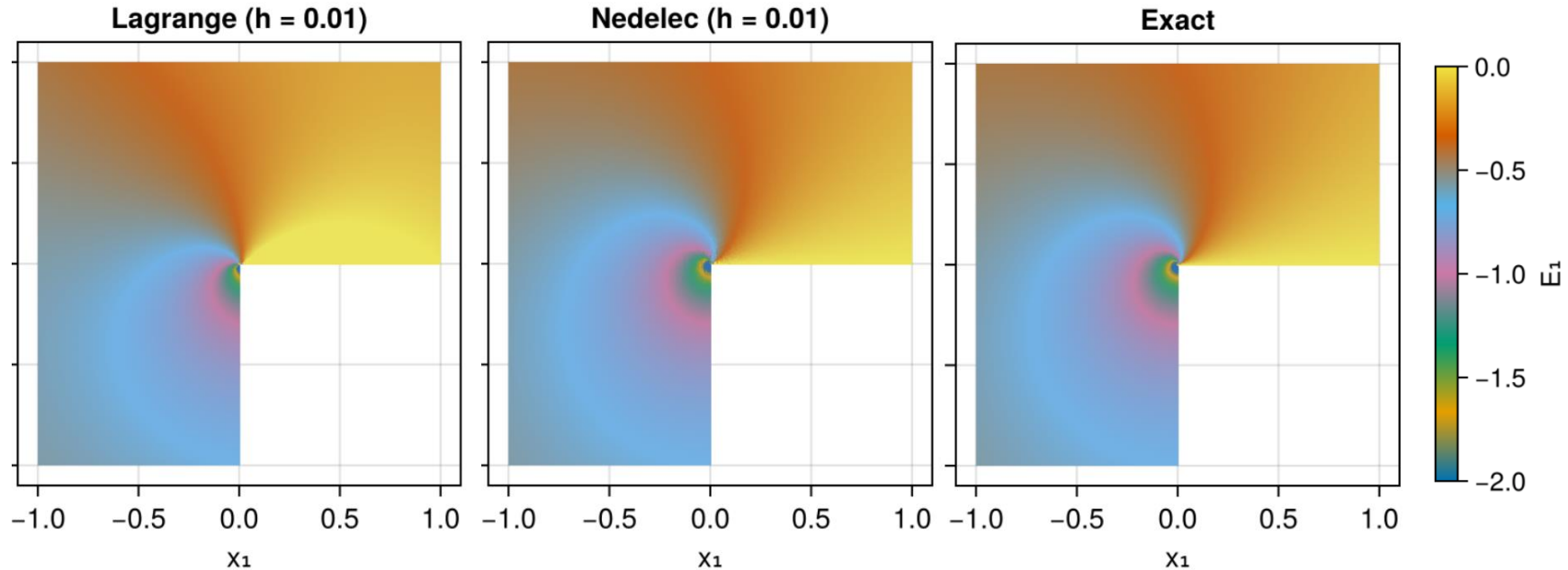
Choose analytical $\mathbf{E}$ and calculate g

$$\mathbf{E}_{\text{exact}}(\mathbf{x}) = \text{grad}(r^{2/3} \sin(2\theta/3))$$



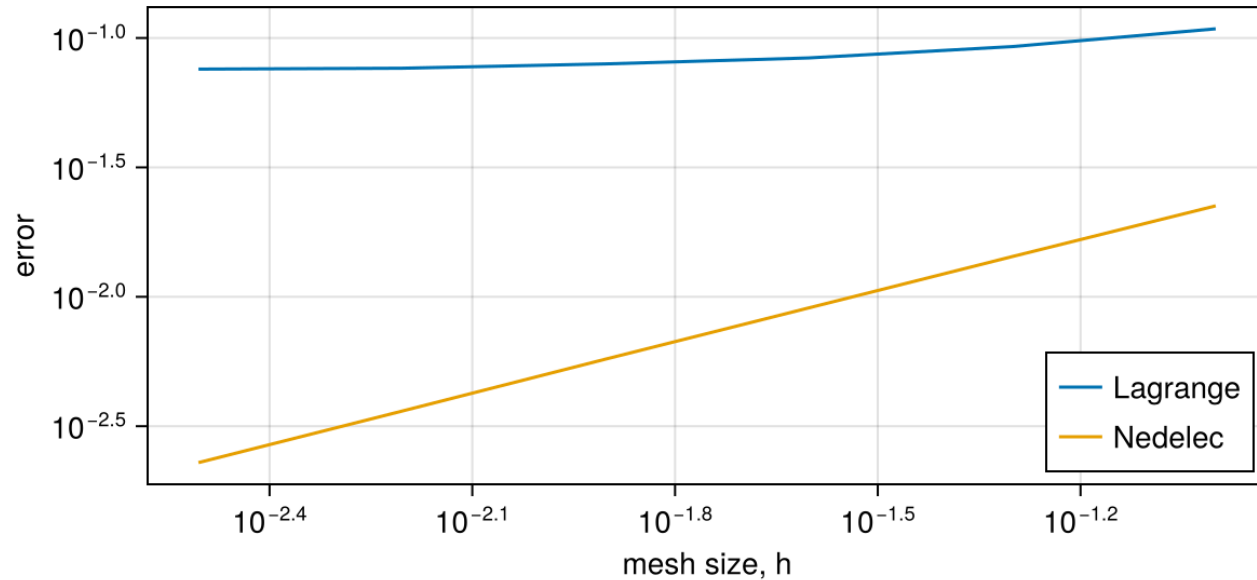
Jay Gopalakrishnan (Portland State University) Maxwell Discretizations: The Good, The Bad & The Ugly
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Ferrite now has true vector interpolations

Conclusions

What to use it for?

- Useful if you work with Maxwell's equations – $H(\text{curl})$
- Applications for dual/mixed formulations with fluxes
- Can be useful for correct postprocessing

What's next?

- Some missing implementations in 3D for $H(\text{curl})$ (issues/prs are open)
- True tensor-valued interpolations (`Double[Contra/Co]variantPiolaMapping`)
- Support in IGA.jl (@Basavesh)
- `CellMultiValues` 😊



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